

FDEs and FPDEs: stability results, diffusive representations and applications

Denis MATIGNON

ISAE, DMIA department & Université de Toulouse.

DFA'10 : *Dynamiques Fractionnaires et Applications*.

— Tuesday, June 1st, 2010 —

1 Motivation

2 BIBO-stability of the Lokshin model

- Analytic solution of a fractional PDE
- Results on Fractional Differential Equations of commensurate orders

3 Asymptotic stability of the Webster-Lokshin model

- Rewriting the model
- Diffusive Pseudo-differential Operators
- Existence and Uniqueness
- Asymptotic stability : a difficult question
- Some numerical illustrations

4 Extensions

- Case of Non-Linear systems
- An open question : Webster-Lokshin FPDE with Bessel-Struve radiation impedance ?
- An open question : perfectly matched impedance for the Euler-Bernoulli beam

Outline

1 Motivation

2 BIBO-stability of the Lokshin model

- Analytic solution of a fractional PDE
- Results on Fractional Differential Equations of commensurate orders

3 Asymptotic stability of the Webster-Lokshin model

- Rewriting the model
- Diffusive Pseudo-differential Operators
- Existence and Uniqueness
- Asymptotic stability : a difficult question
- Some numerical illustrations

4 Extensions

- Case of Non-Linear systems
- An open question : Webster-Lokshin FPDE with Bessel-Struve radiation impedance ?
- An open question : perfectly matched impedance for the

The Lokshin model is non standard

$$\text{Let } (\partial_t^2 + 2\eta \partial_t^{\frac{3}{2}} + \eta^2 \partial_t^1)w - \partial_x^2 w = 0, \quad t > 0, \quad x \in]0, 1[$$

with init. cond. $w(t = 0) = 0$ and $\partial_t w(t = 0) = 0$,
and **controlled** dynamic boundary conditions at $x = 0$;
the system is being **observed** at $x = 1$.

- the damping is modelled by a fractional derivative.

The Lokshin model is non standard

$$\text{Let } (\partial_t^2 + 2\eta \partial_t^{\frac{3}{2}} + \eta^2 \partial_t^1)w - \partial_x^2 w = 0, \quad t > 0, \quad x \in]0, 1[$$

with init. cond. $w(t = 0) = 0$ and $\partial_t w(t = 0) = 0$,
and **controlled** dynamic boundary conditions at $x = 0$;
the system is being **observed** at $x = 1$.

- the damping is modelled by a fractional derivative.
- there is no simple energy property, unlike in the classical cases of *fluid* ($\partial_t^1 w$) or *structural* ($-\partial_t^1 \partial_x^2 w$) dampings.

The Lokshin model is non standard

$$\text{Let } (\partial_t^2 + 2\eta \partial_t^{\frac{3}{2}} + \eta^2 \partial_t^1)w - \partial_x^2 w = 0, \quad t > 0, \quad x \in]0, 1[$$

with init. cond. $w(t = 0) = 0$ and $\partial_t w(t = 0) = 0$,
and **controlled** dynamic boundary conditions at $x = 0$;
the system is being **observed** at $x = 1$.

- the damping is modelled by a fractional derivative.
- there is no simple energy property, unlike in the classical cases of *fluid* ($\partial_t^1 w$) or *structural* ($-\partial_t^1 \partial_x^2 w$) dampings.
- the spatial modes are no more orthogonal.

The Lokshin model is non standard

$$\text{Let } (\partial_t^2 + 2\eta \partial_t^{\frac{3}{2}} + \eta^2 \partial_t)w - \partial_x^2 w = 0, \quad t > 0, \quad x \in]0, 1[$$

with init. cond. $w(t=0) = 0$ and $\partial_t w(t=0) = 0$,
and **controlled** dynamic boundary conditions at $x = 0$;
the system is being **observed** at $x = 1$.

- the damping is modelled by a fractional derivative.
- there is no simple energy property, unlike in the classical cases of *fluid* ($\partial_t^1 w$) or *structural* ($-\partial_t^1 \partial_x^2 w$) dampings.
- the spatial modes are no more orthogonal.
- in the case of the Webster-Lokshin model, the coefficients are variable with space : $\eta \mapsto \eta(x)$.

Outline

1

Motivation

2

BIBO-stability of the Lokshin model

- Analytic solution of a fractional PDE
- Results on Fractional Differential Equations of commensurate orders

3

Asymptotic stability of the Webster-Lokshin model

- Rewriting the model
- Diffusive Pseudo-differential Operators
- Existence and Uniqueness
- Asymptotic stability : a difficult question
- Some numerical illustrations

4

Extensions

- Case of Non-Linear systems
- An open question : Webster-Lokshin FPDE with Bessel-Struve radiation impedance ?
- An open question : perfectly matched impedance for the

Analytic solution of a fractional PDE

Let $(\partial_t^2 + 2\eta \partial_t^{\frac{3}{2}} + \eta^2 \partial_t^1)w - \partial_x^2 w = 0$, $t > 0$, $x \in]0, 1[$

with init. cond. $w(t=0) = 0$ and $\partial_t w(t=0) = 0$, and dynamical boundary conditions of *absorbing* type ($a_0 b_0 > 0$, $a_1 b_1 > 0$) and controlled at $x = 0$:

$$\begin{cases} [a_0 (\partial_t + \eta \partial_t^{\frac{1}{2}}) + b_0 \partial_{-x}] w(t, x=0) = a_0 (\partial_t + \eta \partial_t^{\frac{1}{2}}) u(t) \\ [a_1 (\partial_t + \eta \partial_t^{\frac{1}{2}}) + b_1 \partial_x] w(t, x=1) = 0 \end{cases}$$

with output : $y(t) = w(t, x=1)$. Then $y = h \star u$, with

$$h(t) = \sum_{n=-\infty}^{+\infty} c_n^\eta \left\{ \mathcal{E}_{\frac{1}{2}}(\sigma_n^{\eta+}, t) - \mathcal{E}_{\frac{1}{2}}(\sigma_n^{\eta-}, t) \right\} \in L^1(\mathbb{R}^+) \cap C^\infty(\mathbb{R}^+)$$

where the $\sigma_n^{\eta\pm}$ are the roots of $\sigma^2 + \eta\sigma = s_n^0 = -\alpha^0 + 2i\pi f_n^0$.

BIBO stability comes from $\arg(\sigma_n^{\eta\pm}) > \frac{\pi}{4}$, $\forall n \in \mathbb{Z} \dots$ **Why?**

Fractional Differential Equations

For $0 < \alpha < 1$, consider the input u – output y relation :

$$\sum_{k=0}^p a_k D^{k\alpha} y(t) = \sum_{l=0}^q b_l D^{l\alpha} u(t),$$

It is a *causal* pseudo-differential system, the symbol of which is, by Laplace transf. in some right-half plane $\mathbb{C}_a^+ := \Re e(s) > a$:

$$\mathcal{H}(s) = \frac{Q(s^\alpha)}{P(s^\alpha)} \quad \text{avec} \quad \left\{ \begin{array}{l} Q(\sigma) \triangleq \sum_{l=0}^q b_l \sigma^l \\ P(\sigma) \triangleq \sum_{k=0}^p a_k \sigma^k \end{array} \right. .$$

Necessary and Sufficient Stability Condition

From the input-output viewpoint, the BIBO-stability result is as follows, $y = h \star u$, with :

Theorem

$$\text{BIBO stability} \iff \begin{cases} q \leq p \\ |\arg \sigma| > \alpha \frac{\pi}{2}, \quad \forall \sigma \in \mathbb{C}, / P(\sigma) = 0 \end{cases}$$

In which case, h has the long memory asymptotics :

$$h(t) \sim K t^{-1-\alpha} \quad \text{as } t \rightarrow +\infty.$$

Sketch of the proof

Algebraic tools can be used, since the orders are commensurate to the same α : let P and Q two coprime polynomials, and let $R = Q/P$ the rational function, we get the **structure result** :

Proposition

$$h(t) = \sum_{n=1}^N \sum_{m=1}^{m_n} r_{nm} \mathcal{E}_{\alpha}^{*m}(\lambda_n, t),$$

with $R(\sigma) = \sum_{n=1}^N \sum_{m=1}^{m_n} r_{nm} (\sigma - \lambda_n)^{-m}$.

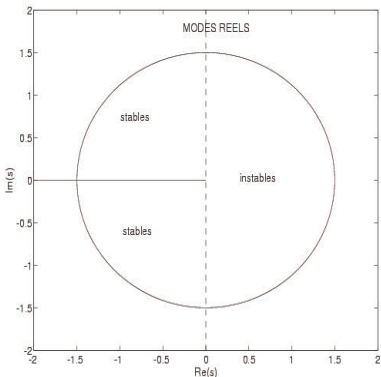
where $\mathcal{E}_{\alpha}^{*m}(\lambda, t)$ is a Mittag-Leffler function (a hypergeometric special function), the LT of which is $(s^{\alpha} - \lambda)^{-m}$.

For $\alpha = 1$, it reduces to the well-known causal polynomial-exponential $\frac{1}{m!} t^{m-1} \exp(\lambda t)$.

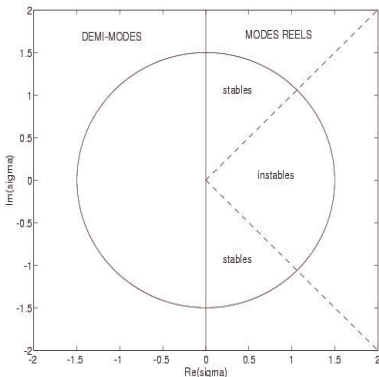
N. & S. Stability condition : an illustration

Stability of $E_\alpha(\lambda t^\alpha)$ with LT $s^{\alpha-1}(s^\alpha - \lambda)^{-1}$, as a fct. of $\arg(\lambda)$.

Laplace plane : s



σ -plane

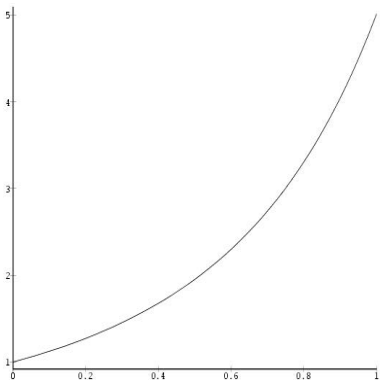


Mittag-Leffler functions in \mathbb{C} (I)

$$t \mapsto E_{\alpha}(\lambda t^{\alpha}) \text{ for } \alpha = \frac{1}{2} \text{ and } \arg(\lambda) = 0$$

Real part

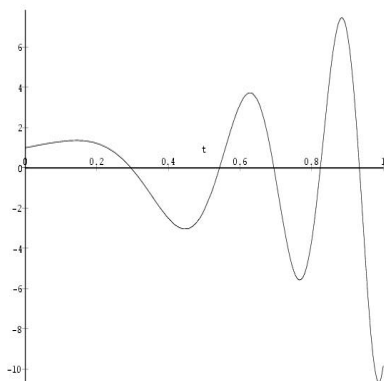
Imaginary part



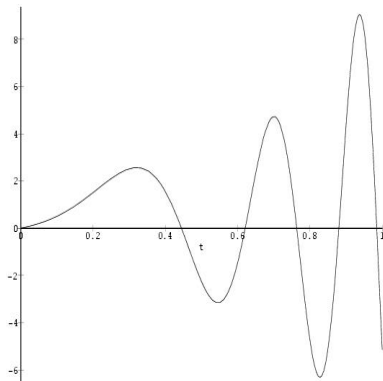
Mittag-Leffler functions in \mathbb{C} (II)

$$t \mapsto E_{\alpha}(\lambda t^{\alpha}) \text{ for } \alpha = \frac{1}{2} \text{ and } \arg(\lambda) = \pi/8$$

Real part



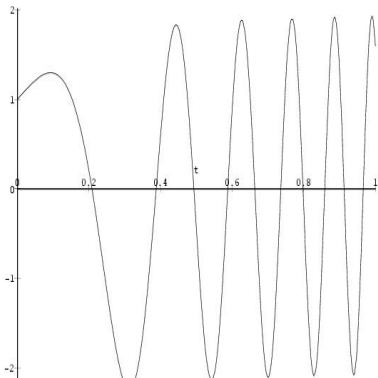
Imaginary part



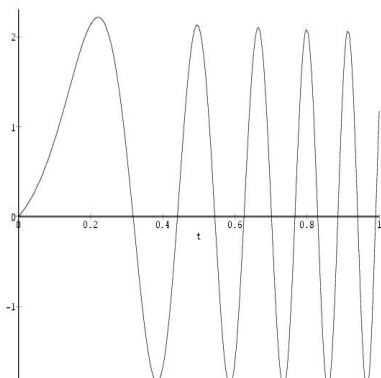
Mittag-Leffler functions in \mathbb{C} (III)

$$t \mapsto E_{\alpha}(\lambda t^{\alpha}) \text{ for } \alpha = \frac{1}{2} \text{ and } \arg(\lambda) = \pi/4$$

Real part



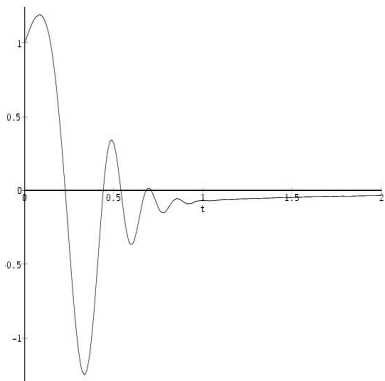
Imaginary part



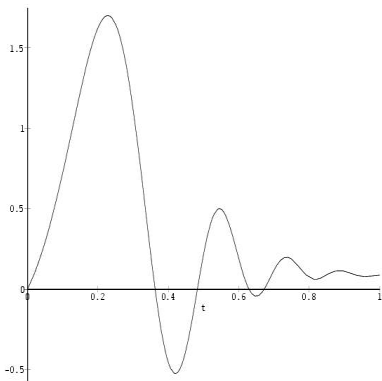
Mittag-Leffler functions in \mathbb{C} (IV)

$$t \mapsto E_\alpha(\lambda t^\alpha) \text{ for } \alpha = \frac{1}{2} \text{ and } \arg(\lambda) = 3\pi/8$$

Real part



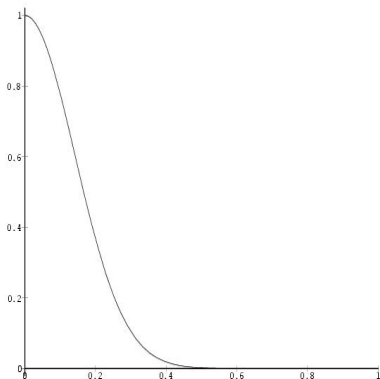
Imaginary part



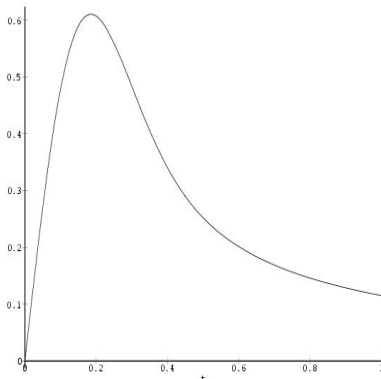
Mittag-Leffler functions in \mathbb{C} (V)

$$t \mapsto E_\alpha(\lambda t^\alpha) \text{ for } \alpha = \frac{1}{2} \text{ and } \arg(\lambda) = \pi/2$$

Real part



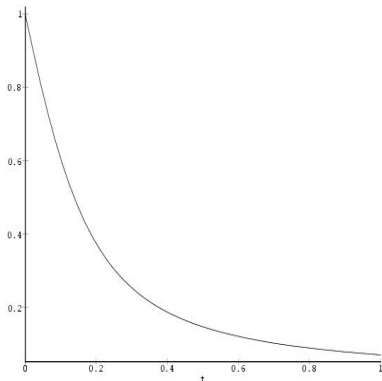
Imaginary part



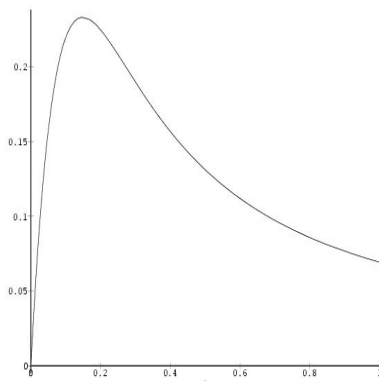
Mittag-Leffler functions in \mathbb{C} (VI)

$$t \mapsto E_{\alpha}(\lambda t^{\alpha}) \text{ for } \alpha = \frac{1}{2} \text{ and } \arg(\lambda) = 3\pi/4$$

Real part



Imaginary part

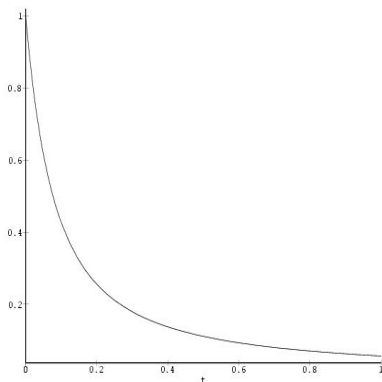


Mittag-Leffler functions in \mathbb{C} (VII)

$$t \mapsto E_{\alpha}(\lambda t^{\alpha}) \text{ for } \alpha = \frac{1}{2} \text{ and } \arg(\lambda) = \pi$$

Real part

Imaginary part



Outline

- 1 **Motivation**
- 2 **BIBO-stability of the Lokshin model**
 - Analytic solution of a fractional PDE
 - Results on Fractional Differential Equations of commensurate orders
- 3 **Asymptotic stability of the Webster-Lokshin model**
 - Rewriting the model
 - Diffusive Pseudo-differential Operators
 - Existence and Uniqueness
 - Asymptotic stability : a difficult question
 - Some numerical illustrations
- 4 **Extensions**
 - Case of Non-Linear systems
 - An open question : Webster-Lokshin FPDE with Bessel-Struve radiation impedance ?
 - An open question : perfectly matched impedance for the

Webster-Lokshin fractional PDE

For $z \in (0, 1)$, with $r(z) > 0$, $\eta(z), \varepsilon(z) \geq 0$, $w(t, z)$ satisfies :

$$\partial_t^2 w + \eta(z) \partial_t^{3/2} w + \varepsilon(z) \partial_t^{1/2} w - \frac{1}{r^2} \partial_z (r^2 \partial_z w) = 0;$$

with static boundary conditions in $z = 0$ and $z = 1$.

- This is equivalent to the first-order system in (p, v) :

$$\begin{aligned} \partial_t p &= -r^{-2} \partial_z v - \varepsilon \partial_t^{-1/2} p - \eta \partial_t^{1/2} p, \\ \partial_t v &= -r^2 \partial_z p, \\ p(z = 0, t) &= 0 \quad \text{and} \quad v(z = 1, t) = 0. \end{aligned}$$

- Use of standard DR for $\partial_t^{-1/2}$, and extended DR for $\partial_t^{1/2}$.

Standard DR : definitions

Let M a positive measure on \mathbb{R}^+ satisfying the **well-posedness** condition (WP) :

$$c_M \triangleq \int_0^{\infty} \frac{dM}{1+\xi} < +\infty.$$

- We define the **dynamical system** with input $u \in L^2(0, T)$, output $y \in L^2(0, T)$ and state $\phi \in H_M = L^2(\mathbb{R}^+, dM)$:

$$\partial_t \phi(\xi, t) = -\xi \phi(\xi, t) + u(t); \quad \phi(\xi, 0) = 0, \quad \forall \xi \in \mathbb{R}^+,$$

$$y(t) = \int_0^{+\infty} \phi(\xi, t) dM(\xi).$$

- Then, $y = h_M \star u$ where the **impulse response** can be written as $h_M(t) = \int_0^{\infty} e^{-\xi t} dM(\xi)$ for $t > 0$.
- The **transfer function** is $\mathcal{H}_M(s) = \int_0^{\infty} \frac{dM(\xi)}{s+\xi}$, in \mathbb{C}_0^+ .

Standard DR : energy balance

The following energy balance is fulfilled, $\forall T > 0$:

$$\int_0^T \mathbf{u}(t) \mathbf{y}(t) dt = \frac{1}{2} \int_0^{+\infty} \phi(\xi, T)^2 dM + \int_0^T \int_0^{+\infty} \xi \phi(\xi, t)^2 dM dt ,$$

where the right-hand side can be decomposed into two parts :

- a *storage* function, evaluated at time T only,

$$E_\phi(T) := \frac{1}{2} \|\phi(T)\|_{H_M}^2,$$
- a residual energy dissipated along the time interval $(0, T)$.

Example : $M_\beta(d\xi) \triangleq \frac{\sin \beta \pi}{\pi} \xi^{-\beta} d\xi$ for $0 < \Re e(\beta) < 1$ fulfills (WP), which gives rise to a *diagonal realization* of the *fractional integral operator* of order β , the transfer function of which is $\mathcal{H}_\beta(s) = s^{-\beta}$.

Note : standard DR belong to the class of *well-posed systems*.

Extended DR : definitions

Let N a positive measure on \mathbb{R}^+ satisfying the **well-posedness** condition (WP).

- We define the **dynamical system** with input $u \in H^1(0, T)$, output $z \in L^2(0, T)$ and state $\tilde{\phi} \in \tilde{H}_N = L^2(\mathbb{R}^+, \xi dN)$:

$$\begin{aligned} \partial_t \tilde{\phi}(\xi, t) &= -\xi \tilde{\phi}(\xi, t) + u(t); & \tilde{\phi}(\xi, 0) &= 0 \quad \forall \xi \in \mathbb{R}^+, \\ z(t) &= \int_0^{+\infty} \partial_t \tilde{\phi}(\xi, t) dN(\xi) = \int_0^{+\infty} \left[u(t) - \xi \tilde{\phi}(\xi, t) \right] dN(\xi). \end{aligned}$$

- Then, $z = \tilde{h}_N \star u = \frac{d}{dt}(h_N \star u)$ with derivative *in the sense of distributions* : the **impulse response** is the distribution $\tilde{h}_N = \frac{d}{dt} \int_0^\infty e^{-\xi t} dN(\xi)$. (N.B. One can also write $z = h_N \star \frac{d}{dt} u$ in the sense of functions).
- The **transfer function** reads $\tilde{H}_N(s) = s \int_0^\infty \frac{dN(\xi)}{s+\xi}$, in \mathbb{C}^+ .

Extended DR : energy balance

The following energy balance is fulfilled, $\forall T > 0$:

$$\int_0^T \mathbf{u}(t) \mathbf{z}(t) dt = \frac{1}{2} \int_0^{+\infty} \xi \tilde{\phi}(\xi, T)^2 dN + \int_0^T \int_0^{+\infty} (\mathbf{u} - \xi \tilde{\phi})^2 dN dt$$

where the right-hand side can be decomposed into two parts :

- a *storage* function, evaluated at time T only,

$$\tilde{E}_{\tilde{\phi}}(T) = \frac{1}{2} \|\tilde{\phi}(T)\|_{H_N}^2,$$

- a residual energy dissipated along the time interval $(0, T)$.

Example : $N_\alpha(d\xi) \triangleq M_{1-\alpha}(d\xi) = \frac{\sin \alpha \pi}{\pi} \xi^{-(1-\alpha)} d\xi$ for

$0 < \Re(\alpha) < 1$ fulfills (WP), which gives rise to a *diagonal realization* of the *fractional derivative operator* of order α , the transfer function of which is $\tilde{\mathcal{H}}_\alpha(s) = s \cdot s^{-(1-\alpha)} = s^\alpha$.

Note : Extended DR **do not** belong to the class of *well-posed systems*, in general.

Existence et uniqueness (I)

With $L_p^2 = \{p, \int_0^1 p^2 r^2 dz < \infty\}$, $L_v^2 = \{v, \int_0^1 v^2 r^{-2} dz < \infty\}$,
 and $\mathcal{H} = L_p^2 \times L_v^2 \times L^2(0, 1; H_M; \varepsilon r^2 dz) \times L^2(0, 1; \tilde{H}_N; \eta r^2 dz)$,
 the system can be put in the abstract form $\partial_t X + \mathcal{A} X = 0$,
 where :

$$\mathcal{A} \begin{pmatrix} p \\ v \\ \varphi \\ \tilde{\varphi} \end{pmatrix} = \begin{pmatrix} r^{-2} \partial_z v + \varepsilon \int_0^{+\infty} \varphi dM + \eta \int_0^{+\infty} [p - \xi \tilde{\varphi}] dN \\ r^2 \partial_z p \\ \xi \varphi - p \\ \xi \tilde{\varphi} - p \end{pmatrix};$$

$$D(\mathcal{A}) = \left\{ (p, v, \varphi, \tilde{\varphi})^T \in \mathcal{V}, \left| \begin{array}{l} p(0) = 0 \\ v(1) = 0 \\ (p - \xi \varphi) \in L^2(0, 1; H_M; \varepsilon r^2 dz) \\ (p - \xi \tilde{\varphi}) \in L^2(0, 1; V_N; \eta r^2 dz) \end{array} \right. \right\}.$$

with $\mathcal{V} = H_p^1 \times H_v^1 \times L^2(0, 1; V_M; \varepsilon r^2 dz) \times L^2(0, 1; \tilde{H}_N; \eta r^2 dz)$.

Existence et uniqueness (II)

Theorem

The operator $\mathcal{A} : D(\mathcal{A}) \subset \mathcal{H} \rightarrow \mathcal{H}$ is maximal monotone.

The *monotonicity* of \mathcal{A} comes from the energy identity :

$\forall X \in D(\mathcal{A}),$

$$(\mathcal{A}X, X)_{\mathcal{H}} = \int_0^1 \|\varphi\|_{H_M}^2 \varepsilon r^2 dz + \int_0^1 \|p - \xi \tilde{\varphi}\|_{H_N}^2 \eta r^2 dz \geq 0 .$$

Corollary

Hille-Yosida theorem enables to conclude to the existence and uniqueness of a strong solution for the original problem.

Note : in case of dynamical boundary conditions, the Kalman-Yakubovich-Popov lemma will be used to realize the output impedance, which is a positive real rational function of s .

Internal asymptotic stability

The proof is difficult. Why ?

- In fact, LaSalle's invariance principle requires, in infinite dimension, the hypothesis of precompactnes of the trajectories ; but this latter hypothesis cannot be checked a priori for diffusive realizations, since a diffusion equation in an unbounded domain is hidden behind them, and the canonical injection from $H^1(\mathbb{R})$ into $L^2(\mathbb{R})$ is not compact (Rellich theorem does not apply).

Internal asymptotic stability

The proof is difficult. **Why ?**

- In fact, **LaSalle's invariance principle** requires, in infinite dimension, the hypothesis of **precompactnes** of the trajectories ; but this latter hypothesis cannot be checked a priori for diffusive realizations, since a diffusion equation in an **unbounded domain** is hidden behind them, and the canonical injection from $H^1(\mathbb{R})$ into $L^2(\mathbb{R})$ is not compact (Rellich theorem does not apply).
- The refined spectral analysis of the infinitesimal generator $-\mathcal{A}$ of the semigroup on the Hilbert state \mathcal{H} enables to use the **stability result by Arendt–Batty** or Lyubich–Phong, et helps prove the result of *internal asymptotic stability* ; the proof is quite involved (Lax–Milgram theorem for the FDE, and Fredholm alternative for the FPDE).

Definition and Analysis of numerical schemes

- Fractional derivatives are difficult to numerically approximate, and usually involve **hereditary algorithms**, thus turning into memory storage problems on the computer.

Definition and Analysis of numerical schemes

- Fractional derivatives are difficult to numerically approximate, and usually involve **hereditary algorithms**, thus turning into memory storage problems on the computer.
- Standard numerical approximations of the extended system with DR enable to **define** memoryless numerical schemes ; more precisely the schemes have **finite memory**, once the $\{\xi_j\}_{1 \leq j \leq J}$ have been chosen.

Definition and Analysis of numerical schemes

- Fractional derivatives are difficult to numerically approximate, and usually involve **hereditary algorithms**, thus turning into memory storage problems on the computer.
- Standard numerical approximations of the extended system with DR enable to **define** memoryless numerical schemes ; more precisely the schemes have **finite memory**, once the $\{\xi_j\}_{1 \leq j \leq J}$ have been chosen.
- The proof of **convergence** of the numerical schemes is based on **discrete extended energy** techniques, which mimic the principle of the extended energy for the continuous system.

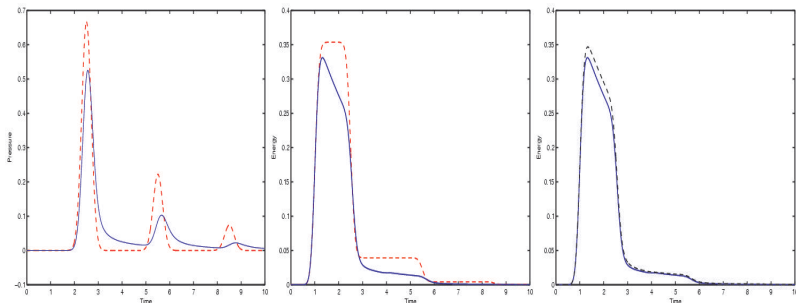
Definition and Analysis of numerical schemes

- Fractional derivatives are difficult to numerically approximate, and usually involve **hereditary algorithms**, thus turning into memory storage problems on the computer.
- Standard numerical approximations of the extended system with DR enable to **define** memoryless numerical schemes ; more precisely the schemes have **finite memory**, once the $\{\xi_j\}_{1 \leq j \leq J}$ have been chosen.
- The proof of **convergence** of the numerical schemes is based on **discrete extended energy** techniques, which mimic the principle of the extended energy for the continuous system.
- Some **Illustrations** !

Some numerical illustrations

Influence of parameter η (I)

- Output signal. Wave & augmented (- -) energies

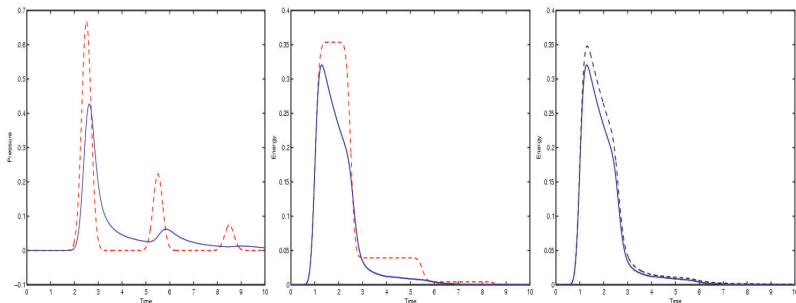


- For a cylinder, in blue $\eta = 0.1$, in red $\eta = 0$.

Some numerical illustrations

Influence of parameter η (II)

- Output signal. Wave & augmented (- -) energies.

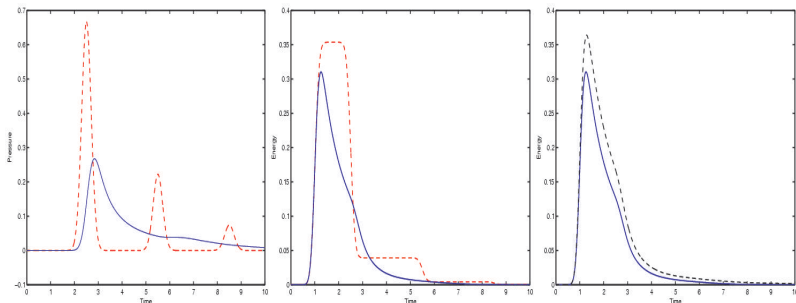


- For a cylinder, in blue $\eta = 0.2$, in red $\eta = 0$.

Some numerical illustrations

Influence of parameter η (III)

- Output signal. Wave & augmented (- -) energies.

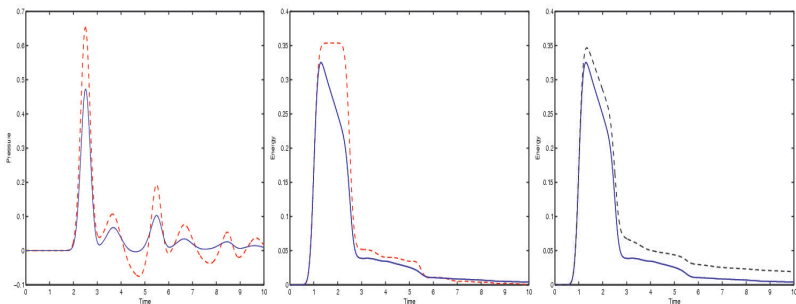


- For a cylinder, in blue $\eta = 0.5$, in red $\eta = 0$.

Some numerical illustrations

Example of a trapped modes

- Output signal. Wave & augmented (- -) energies.



- Trapped modes in a duct with two cones facing each other, in blue $\varepsilon = 0.2$ and $\eta = 0.05$, in red $\varepsilon = \eta = 0$.

Outline

- 1 **Motivation**
- 2 **BIBO-stability of the Lokshin model**
 - Analytic solution of a fractional PDE
 - Results on Fractional Differential Equations of commensurate orders
- 3 **Asymptotic stability of the Webster-Lokshin model**
 - Rewriting the model
 - Diffusive Pseudo-differential Operators
 - Existence and Uniqueness
 - Asymptotic stability : a difficult question
 - Some numerical illustrations
- 4 **Extensions**
 - Case of Non-Linear systems
 - An open question : Webster-Lokshin FPDE with Bessel-Struve radiation impedance ?
 - An open question : perfectly matched impedance for the

Un résultat général

Theorem

Soit H un espace de Hilbert, soit $A : D(A) \subset H \rightarrow H$ un **opérateur maximal monotone**, et F une fonction **non-linéaire** $F : H \rightarrow H$ telle que le pb. d'évolution semi-linéaire :

$$\partial_t X + AX = F(X), \quad \text{et } X(0) = X_0 \in D(A)$$

soit **bien posé**, pour $t \in [0, T_{max})$, au sens de l'existence et de l'unicité de $X \in C^1([0, T_{max}); H) \cap C^0([0, T_{max}); D(A))$, une solution forte.

Alors, pour deux OPD de type diffusif et **positif**, l'un standard h_{M^*} et l'autre étendu par dérivation \widetilde{h}_{N^*} , le pb. pseudo-différentiel non-linéaire :

$$\partial_t X + h_{M^*} X + \widetilde{h}_{N^*} X + AX = F(X), \quad \text{et } X(0) = X_0 \in D(A)$$

est **bien posé**, pour $t \in [0, T'_{max})$, au sens d'une unique solution forte $X \in C^1([0, T'_{max}); H) \cap C^0([0, T'_{max}); D(A))$.

Perturbation diffusive de systèmes non-linéaires conservatifs

Remark

L'hypothèse du théorème précédent est vérifiée dans le cas où la non-linéarité est **localement lipschitzienne** sur H .

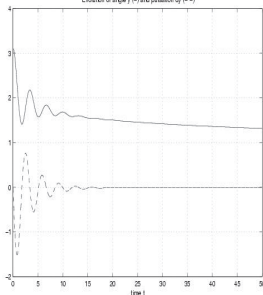
Corollary

*Si le système différentiel non-linéaire de départ est **conservatif**, alors comme l'énergie étendue associée au système perturbé est décroissante et bornée par sa valeur initiale, le résultat d'existence locale se prolonge en un résultat d'existence **globale**, i.e. $T'_{max} = +\infty$.*

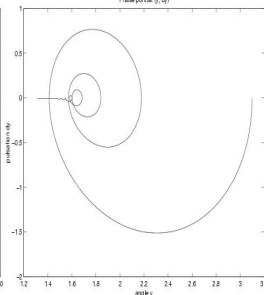
Résultats de simulation (I) : pendule **linéaire**

- $\ddot{\vartheta} + \eta \partial_t^{-\beta} \dot{\vartheta} + \vartheta = 0$, pour $\beta = 0.75$ et $(\vartheta_0, \dot{\vartheta}_0) = (3.5, 0)$.

(G)

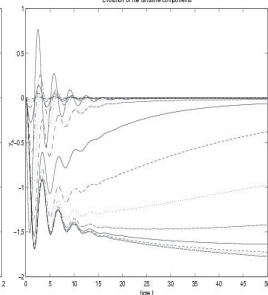
Evolution of angle ϑ (-) and pulsation $\dot{\vartheta}$ (- -)

(C)

Phase portrait ($\vartheta, \dot{\vartheta}$)

(D)

Evolution of the diffusive components



- (G) : Évolution de l'angle ϑ et de la vitesse angulaire $\dot{\vartheta}$,
- (C) : Section du portrait de phase dans le plan $(\vartheta, \dot{\vartheta})$,
- (D) : Évolution des composantes diffuses $\{\phi_k(t) = \phi(t, \xi_k)\}_{1 \leq k \leq K}$ pour $K = 25$.

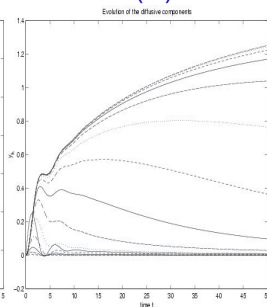
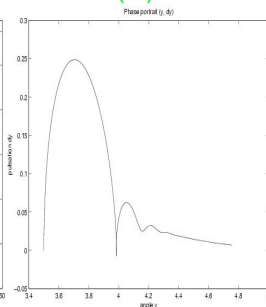
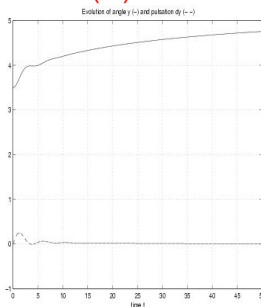
Résultats de simulation (II) : pendule **non**-linéaire

- $\ddot{\vartheta} + \eta \partial_t^{-\beta} \dot{\vartheta} + \sin(\vartheta) = 0$; $\beta = 0.75$ et $(\vartheta_0, \dot{\vartheta}_0) = (3.5, 0)$.

(G)

(C)

(D)

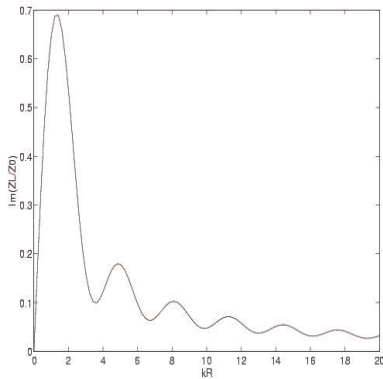
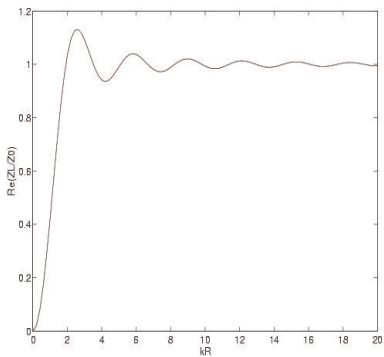


- (G) : Évolution de l'angle ϑ et de la vitesse angulaire $\dot{\vartheta}$,
- (C) : Section du portrait de phase dans le plan $(\vartheta, \dot{\vartheta})$,
- (D) : Évolution des composantes diffusives $\{\phi_k(t) = \phi(t, \xi_k)\}_{1 \leq k \leq K}$ pour $K = 25$.

An open question : Webster-Lokshin FPDE with Bessel-Struve radiation impedance ?

Is the following radiation impedance diffusive of the second kind ?

Bessel J_1 and Struve H_1 special functions



Real p. $x = kR \mapsto 1 - \frac{J_1(2x)}{2x}$

Imaginary p. $x = kR \mapsto \frac{H_1(2x)}{x}$

An open question : perfectly matched impedance for the Euler-Bernoulli beam

Matrix-valued impedances as diffusive PDOs ?

As an example, in order to solve the **impedance matching** problem for the Euler–Bernoulli beam, an impedance matrix of the form

$$\mathcal{Z} = \begin{bmatrix} a\partial_t^{+\alpha} & 1 \\ 1 & b\partial_t^{-\alpha} \end{bmatrix}$$






is found in the time domain, with $\alpha = \frac{1}{2}$.

A **necessary and sufficient condition** for the positivity of this operator is :





$$ab > \frac{1}{(\cos \alpha\pi)^2} .$$

Open question : is there any dissipative diffusive realization of the transfer matrix $\hat{\mathcal{Z}}(s)$, which is positive in the sense $\forall s \in \mathbb{C}_0^+$, $\hat{\mathcal{Z}}(s) + \hat{\mathcal{Z}}(s)^H \geq 0$, i.e. a positive symmetric real-valued matrix ?

Some references

-  W. ARENDT AND C. J. K. BATTY, *Tauberian theorems and stability of one-parameter semigroups*, Trans. Amer. Math. Soc., 306 (1988), pp. 837–852.
-  Y. LYUBICH AND V. PHÓNG, *Asymptotic stability of linear differential equations on Banach spaces*, Studia Mathematica, 88 (1988), pp. 37–42.
-  O. J. STAFFANS, *Well-posedness and stabilizability of a viscoelastic equation in energy space*, Trans. Amer. Math. Soc., 345 (1994), pp. 527–575.
-  D. MATIGNON, *Stability properties for generalized fractional differential systems*, ESAIM : Proc., 5 (1998), pp. 145–158.
-  H. ZWART, *Transfer functions for infinite-dimensional systems*, Systems Control Lett., 52 (2004), pp. 247–255.

Some more references

-  H. HADDAR, T. HÉLIE, AND D. MATIGNON, *A Webster-Lokshin model for waves with viscothermal losses and impedance boundary conditions : strong solutions*, in WAVES'03, Jyväskylä, Finland, jul. 2003, INRIA, pp. 66–71.
-  D. MATIGNON AND H. ZWART, *Standard diffusive systems are well-posed linear systems*, in Mathematical Theory of Networks and Systems (MTNS), Leuven, Belgium, jul 2004.
-  D. MATIGNON AND C. PRIEUR, *Asymptotic stability of linear conservative systems when coupled with diffusive systems*, ESAIM : COCV, 11 (2005), pp. 487–507.
-  T. HÉLIE, D. MATIGNON, *Diffusive representations for the analysis and simulation of flared acoustic pipes with visco-thermal losses*, Mathematical Models and Methods in Applied Sciences, 16 (2006), pp. 503–536.